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Dielectric properties of oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)$ $O_{12-\delta}$ bulk superconductor

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Abstract

Oxygen post-annealing experiments are carried out to optimize the oxygen contents in $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ (y=0, 0.25, 0.5, and 0.75) superconductor samples and to investigate its effects on their dielectric properties {i.e. dielectric constants (ϵ' , ϵ'' , ϵ'' ,), absolute dielectric loss I tan δ I and ac-conductivity (σ_{ac})}. The negative capacitance (NC) is observed in all post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples. Almost all the dielectric properties are suppressed except dielectric loss ($\tan \delta$), which is increased after oxygen post-annealing. The appreciable change in the dielectric properties is the indication of oxygen intake by $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples after oxygen post-annealing. The thermal agitation is higher and polarizability is lower at normal state (300 K), resulting in decrease of the dielectric constant and increase in $\tan \delta$ which is vice versa at lower temperatures. The observed variation in the dielectric behavior of $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples strongly depends on the working (or operating) temperature and frequency of the external applied ac-field. The dielectric behavior of the material can be regulated by the post-annealing in oxygen as a result of which the variation of oxygen content of the materials structures may take place thus changing the carrier's density. On the other hand, this post-annealing may affect the inter-grain connectivity of the materials that can influence the dielectric properties of the material.

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 $Keywords: Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ superconductor; Negative capacitance; Dielectric properties; Post-annealing

1. Introduction

The oxygen content in the charge reservoir layers control the carriers density in CuO_2 planes of multilayered cuprates [1,2]. Oxygen post-annealing is the simple and appropriate way to vary the oxygen contents in the cuprates. The change in O_δ oxygen variation in $Cu_{0.5}Tl_{0.5}Ba_2O_{4-\delta}$ charge reservoir layer controls the flow of carriers toward the conducting CuO_2 planes of $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples. The charge carriers are mobile in the conducting CuO_2 planes, while Tl^{3+} , Cu^{2+} and Ba^{2+} ions localize the charge carriers in $Cu_{0.5}Tl_{0.5}Ba_2O_{4-\delta}$ charge reservoir layer [3]. The dielectric

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constant arises from the accumulation of carriers in between CuO_2 planes and $\text{Cu}_{0.5}\text{Tl}_{0.5}\text{Ba}_2\text{O}_{4-\delta}$ charge reservoir layer. It is well known that there are four primary mechanisms of polarization in materials [3–7] namely electronic polarization (α_e) , which can be observed at very high frequencies of the order of 10^{15} Hz (i.e. in ultraviolet optical range), atomic and ionic polarization (α_a) , which takes place at the frequencies in range of 10^{10} – 10^{13} Hz (i.e. in the infrared optical range), dipolar or oriental polarization (α_o) , which occurs in the range of frequencies from 10^3 to 10^6 Hz (i.e. in the sub-infrared optical range), and interfacial polarization (α_i) , which is more sensitive in the low frequency range of 10^3 Hz and may extend to few kilohertz range [8]. Each mechanism of polarization involves a short-range motion of charges and contributes to the total polarization of the material.

The dipolar or oriental polarization and interfacial polarization mainly contribute to the dielectric response of

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 $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ superconductor samples because the dielectric constant become saturated in the high frequency range at all temperatures from superconducting state to room temperature. The dipolar polarization or the oriental polarization contributes to the dielectric properties in the subinfrared range of frequency [3]. Mobile carriers are impeded by a physical barrier in the interfacial polarization that inhibits charge migration. The charges pile up at the barrier producing a localized polarization in the materials [8]. Therefore, all the above mentioned mechanisms of polarization are important at low frequencies however, the effects of dipolar polarization are weak and the dielectric constant mainly originates from electric and lattice polarization contribution on account of their short relaxation time. The conduction mechanisms of variety of electronic devices strongly depend upon the temperature, frequency, fabrication conditions and process, surface charges, doping concentration and impurities (interface states or dislocations). The dielectric properties do not reflect the same intrinsic response at all frequencies and temperatures. Therefore, the study of dielectric properties and ac-conductivity over the wide range of frequency and temperature is very important. The determination of dielectric properties and ac-conductivity of the materials at different frequencies and temperatures can provide valuable information about the conduction and polarization mechanisms. The materials with giant dielectric constant are playing an increasingly significant role in the growth of microelectronics because of the desire for smaller and more robust devices such as capacitor and memory devices [9]. However, the main problem for potential applications is the high conductivity in cuprates, which results in high dissipation loss. The large polarizability of Tl³⁺, Ba²⁺ and Cu²⁺ ions may contribute to the possible high frequency practical applications of Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-ν}Cd_ν)O_{12-δ} compound. Most importantly, the conductivity can be adjusted in Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-v}Cd_v) $O_{12-\delta}$ compound to relatively low value by changing the oxygen contents because the oxygen contents in the charge reservoir layer determines the valence state of Tl and Cu that control the carriers density in the conducting CuO₂ planes [10-14]. As is known, very low conductivity is absolutely necessary for low dissipation factor.

In this article, we have reported the frequency dependent dielectric measurements of oxygen post-annealed Cu_{0.5}Tl_{0.5}- $Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ superconductor in comparison with as-prepared samples. The oxygen post-annealing experiments are normally carried out to modify the chemical state of elements in the unit cell by varying the oxygen contents to optimize the carriers concentration, to increase the grains size and to enhance the inter-grain connectivity in bulk materials [10-14]. Therefore our main objective of oxygen postannealing experiments on Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-v}Cd_v)O_{12-δ} superconductor samples was to optimize the carriers concentration and to increase the grains size as well as to enhance the inter-grain connectivity. The effects of all these factors after oxygen post-annealing can be visualized in the form of variation in the dielectric properties deduced from experimental data of capacitance and conductance measurements at different temperatures and frequencies. We have calculated the dielectric

constants (ε'_n , ε''_r), absolute dielectric loss Itan δ I and acconductivity (σ_{ac}) by experimentally measuring the capacitance (C) and conductance (G) with the help of Hewlett-Packard 4275 A Multi-Frequency LCR Meter from 10 KHz to 10 MHz at various temperatures from 80–300 K. We have compared the dielectric properties of post-annealed Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-y}Cd_y) O_{12- δ} samples with those of un-annealed as-prepared Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-y}Cd_y)O_{12- δ} samples [15]. The capacitance, conductance, dielectric constant, dielectric loss and ac-conductivity are the main parameters for the selection of the materials for device applications. The dielectric parameters (ε'_n , ε''_r , Itan δ I and σ_{ac}) can be determined by means of C and G measurements in the test frequency (f) range of 10 KHz to 10 MHz using the following expressions [16].

The real part of dielectric constant (ε'_r) can be calculated by the expression

$$\varepsilon'_r = \frac{Cd}{A\varepsilon_o} \tag{1}$$

where *C* is capacitance (*F*), *d* is the thickness of the dielectric/separation between the plates of the capacitor (m), ε_o is the permittivity of free space ($\varepsilon_o = 8.85 \times 10^{-12} \,\mathrm{F m}^{-1}$) and *A* is the area of the electrodes (m²).

The imaginary part of dielectric constant or dielectric loss (ε''_r) of the material can be calculated from the conductance (G) by using the following equation

$$\varepsilon''_{r} = \frac{Gd}{A\varepsilon_{o}\omega} \tag{2}$$

where ω (=2 πf) is the angular frequency.

The dielectric loss $\tan \delta$ can be calculated from the relation

$$\tan \delta = \frac{\varepsilon''_r}{\varepsilon'_r} \tag{3}$$

The ac-conductivity (σ_{ac}) can be calculated by using the following relation

$$\sigma_{ac} = 2\pi f \varepsilon_o \varepsilon_r' \tan \delta \tag{4}$$

where f is the frequency of the external applied ac-field (Hz).

2. Experimental

The polycrystalline bulk Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-v}Cd_v)O_{12-δ} (y=0, 0.25, 0.5, and 0.75) superconductor samples were synthesized by the solid-state reaction method using Cd (NO₃)₂, Ba(NO₃)₂, Ca(NO₃)₂ and Cu(CN) as starting compounds. These compounds are mixed in appropriate ratios and ground in a quartz mortar pestle for about an hour and then fired at 880 °C in a quartz boat for 24 h followed by furnace cooling to room temperature. This material was again ground for about 1 h and heat-treated at the same temperature of 880 °C for 24 h. The fired material was then mixed with Tl₂O₃ and ground again for about 1 h. The material mixed with Tl₂O₃ was palletized under 5 t/cm² pressure and the pellets were enclosed in gold capsules for heat-treatment at 880 °C for 10 min followed by quenching to room temperature. The bar shaped samples of dimensions $2 \text{ mm} \times 2.5 \text{ mm} \times 10 \text{ mm}$ were used for characterizations. The post-annealing experiments on

these samples were carried out in tubular furnace at 500 °C in flowing oxygen for 6 h. The bar shaped samples of dimensions $2 \text{ mm} \times 2.5 \text{ mm} \times 10 \text{ mm}$ were used for dc-resistivity ρ (Ω cm) measurements by the four-probe method. The frequency dependent dielectric measurements were carried out with the help of Hewlett-Packard 4275 A Multi-Frequency LCR Meter in the frequency range of 10 KHz to 10 MHz. The calibrations of the LCR meter were performed with the help of standard capacitor in order to elucidate the influence of an external circuit. The parallel mode of LCR meter was used for these measurements. A conventional two-probe technique was used for the capacitance (C) and conductance (C) measurements of $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ superconductor samples. Silver paint was applied on both the surfaces of the sample and copper leads were fixed on the silver electrode surfaces.

3. Results and discussion

The summary of the comparison of superconductivity and dielectric properties of as-prepared [15] and oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ superconductor samples is given in Table 1.

The x-ray diffraction (XRD) scans of as-prepared Cu_{0.5}Tl_{0.5}- $Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ (y=0, 0.25, and 0.5) samples have shown tetragonal structure following P4/mmm space group with a-axes lengths 4.146, 4.146, and 4.147 Å and c-axes lengths 17.86, 17.83, and 17.8 Å for y=0, 0.25, and 0.5, respectively [17]. The lattice parameters of these samples are calculated by using computer software (crystal). The ionic size of Cd^{2+} (1.41 Å) is large as compared to that of Cu^{2+} (1.17 Å) but the decrease in c-axis length gives the clue of reduced Jahn-Teller distortion with Cd-doping. The non-linear molecules having a degenerate electronic state distort to lower the symmetry of the molecule and reduce the degeneracy. This is commonly known as Jahn-Teller effect and this effect can be observed for partially filled energy level [17]. The dcresistivity ρ (Ω cm) versus temperature T (K) measurements of oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ superconductor samples are shown in Fig. 1. The variation in resistivity for all these samples from room temperature down to onset of superconductivity is metallic with onset temperature (T_c^{onset}) of superconductivity around 131, 124, 120, and 108 K and zero resistivity critical temperature $\{T_c\}$ (R=0)} around 120, 114, 115, and 85 K for y=0, 0.25, 0.5, and 0.75 respectively. Both the T_c^{onset} and T_c (R=0) are shifted to higher temperature values in all the oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ samples. The most plausible reason for the improvement of these superconductivity parameters is the optimal carriers concentration and increased superconducting volume fraction of the samples after oxygen post-annealing.

The variation in the real part of dielectric constant (ε'_r) of oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples as a function of frequency at various temperatures in the range of 80–300 K is shown in Fig. 2(a–d). The portion of energy stored within the material of the sample is given by ε'_r , when exposed to electric field. The negative value of ε'_r is

 \leq 2

$Cu_{0.5}Tl_{0.5}Ba_2Ca_3$ $(Cu_{4-y}Cd_y)O_{12-\delta}$	$Cu_{0.5}Tl_{0.5}Ba_2Ca_3$ T_c $(R=0)$ (K) $(Cu_{4-y}Cd_y)O_{12-\delta}$		$\epsilon_{\rm u,s} T l_{\rm 0.5} B a_{\rm 2} C a_{\rm 3}$ $T_{\rm c} (R=0)$ (K) $\epsilon_{\rm r'}$ at 10 kHz $\epsilon_{\rm r'}$		e", at 10 kHz		tan 8	•	$\sigma_{ac} \; (\Omega \; \mathrm{cm})^{-1}$	
y	As-prepared Oxygen	Oxygen As-prepar post-annealed 80–300 K	As-prepared 80–300 K	Oxygen post-annealed 80–300 K	As-prepared 80–300 K	Oxygen post-annealed 80-300 K	As-prepared Oxygen 80–300 K post-anr 80–300 I	Oxygen post-annealed 80–300 K	As-prepared Oxygen 80-300 K post-anr 80-300 I	Oxygen post-anneak 80–300 K
0	109	120	$-3.7 \times 10^6 \text{ to } -2.6 \times 10^5$	-2.06×10^6	$1.04 \times 10^9 - 2.23 \times 10^8$	$1.04 \times 10^{9} - 2.23 \times 10^{8} 2.39 \times 10^{8} - 1.79 \times 10^{7} 300 - 1205$	300-1205	265–1323	577-124	133–10
0.25	107	114	$-4.9 \times 10^5 \text{ to } -2.8 \times 10^4$	-3.52×10^5 -3.52 × 10^5	$1.5\times 10^{8} - 2.7\times 10^{7}$	$1.63 \times 10^{8} - 2.50 \times 10^{7}$ 286-894	286–894	223–1449	83–15	60–11
0.5	102	115	$-3.5 \times 10^5 \text{ to } -2.7 \times 10^4$	-3.38×10^5	$1.4 \times 10^8 - 1.9 \times 10^7$	$3.30 \times 10^{8} - 2.00 \times 10^{7}$ 258–786	258–786	165–1802	50-12	42–14
0.75	82	85	to -4.27×10^{-5} -6.6×10^{5} to -5.02×10^{4} -3.49×10^{4}	10 -4.27×10^{2} -3.49×10^{4}	$8.9 \times 10^{7} - 2.2 \times 10^{7}$	$1.43 \times 10^8 - 1.91 \times 10^7$ 32-684	32–684	728–2260	40–11	28-11
				to -1.33×10^3						

calculated from the negative capacitance (NC) and the term NC reflects a decrease in sample capacitance from the geometric capacitance without the sample. This is most probably due to the development of positive space charge in the neighborhood of the electrodes. The positive space charge is most likely accumulated at the outer surface of the device

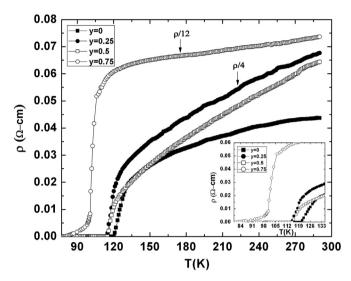


Fig. 1. The resistivity versus temperature measurements of oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ (y=0, 0.25, 0.5, and 0.75) super-conductor samples.

because the free negative charge carriers (electrons) flow toward the metal electrodes from the ceramic sample. As the Fermi level of the ceramics is at higher energy than that of metals, therefore, the migration of electrons from ceramics to metal surface is possible; the ceramics have less filled states than metals. The NC was observed in the samples with metalsemiconductor interfaces in the low frequency ranges [18-23]. The NC was also observed in some other materials [24–26]. The reasons of NC mechanism are different in different materials. In these studies, it is reported that the microscopic phenomena associated with NC at the metal-semiconductor interface includes the trapping of carriers (capture and emission), space charge effects and contact injection [18,27]. The maximum values of ε'_r in oxygen post-annealed Cu_{0.5}Tl_{0.5}- $Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ samples at frequency of 10 KHz and in temperature range of 80-300 K are around from -2.06×10^6 to -1.03×10^4 , -3.52×10^5 to -2.34×10^4 , -3.38×10^5 to -4.27×10^3 and -3.49×10^4 to -1.33×10^3 for y=0, 0.25, 0.5, and 0.75 respectively. The maximum values of ε'_r in all $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ samples is decreased after oxygen post-annealing, Fig. 2(a-d).

The imaginary part of the dielectric constant (ε''_r) gives the absorption and attenuation of energy across the interfaces under the applied external electric field. The interfaces include grain boundaries, localized defects and localized charge densities at the defects sites. The variation of ε''_r versus $\text{Log}\{f(\text{Hz})\}$ at various temperatures from 80–300 K of oxygen

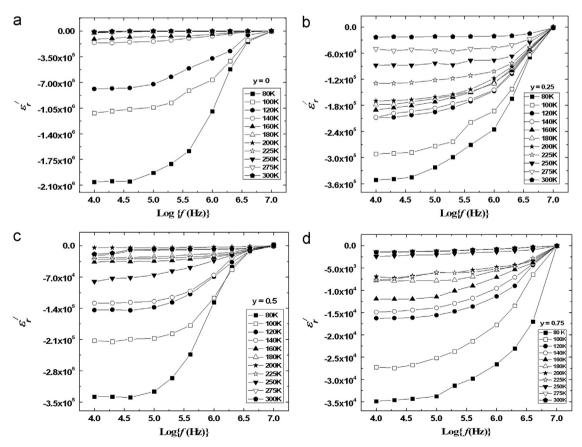


Fig. 2. (a–d). The dielectric constant (ϵ_r') versus log of the frequency $\{\text{Log }f(\text{Hz})\}$ of the external field applied on oxygen post-annealed $\text{Cu}_{0.5}\text{Tl}_{0.5}\text{Ba}_2\text{Ca}_3$ $(\text{Cu}_{4-\nu}\text{Cd}_{\nu})\text{O}_{12-\delta}$ (a) y=0, (b) y=0.25, (c) y=0.5, and (d) y=0.75, superconductor samples.

post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples, Fig. 3(a–d). The maximum values of ε''_r are around from 2.39×10^8 to 1.79×10^7 , 1.63×10^8 to 2.50×10^7 , 3.30×10^8 to 2.00×10^7 and 1.43×10^8 to 1.91×10^7 at 80–300 K and 10 KHz for y=0, 0.25, 0.5, and 0.75, respectively, after oxygen post-annealing. The less attenuation in energy is observed after oxygen post-annealing that may be due to large intake of oxygen in the $Cu_{0.5}Tl_{0.5}Ba_2O_{4-\delta}$ charger reservoir layer. The large intake of oxygen may change the charge state of Tl from (3+) to (1+) in $Cu_{0.5}Tl_{0.5}Ba_2O_{4-\delta}$ charger reservoir layer that reduces the localization of carriers resulting in the decrease in the dielectric properties after oxygen post-annealing [28,29].

The values of ε'_r and ε''_r in $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples are much less than those observed in $Tl_2Ba_2Ca_2Cu_3O_x$ (Tl-2223) and $Tl_2Ba_2Ca_1Cu_2O_x$ (Tl-2212) superconductor samples [6]. The values of the dielectric constants observed in $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples are even much smaller than those values observed in Nd-doped $Bi_2Ba_2Nd_{1.6}Ce_{0.4}$ - $Cu_3O_{10-\delta}$ and Bi-Pb-Sr-Ca-Cu-O superconductors [8,30]. A significant difference in the values of ε'_r and ε''_r may possibly be due to larger difference in the thickness of charge reservoir layer; $Tl_2Ba_2O_{5-\delta}$ charge reservoir layer of double Tl layer compounds is thicker as compared to $TlBa_2O_{4-\delta}$ charge reservoir layer of single Tl layer compounds. Moreover, $(Cu_{0.5}Tl_{0.5})Ba_2O_4$ - δ charge reservoir layer in $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$

superconductor is more conducting, whereas, Tl₂Ba₂O_{5-δ} charge reservoir layer in Tl-2223 and Tl-2212 superconductors is insulating [7]. The former supplies more carriers to the conducting CuO₂/CdO₂ planes as compared to the totally insulating Tl₂Ba₂O_{5-δ} charge reservoir layers. Naturally lower values of capacitance are observed in Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-v}Cd_v)O_{12-δ} samples as compared to Tl-2223 and Tl-2212 samples, whereas a higher conductance is observed in our samples as compared to double layered thallium based superconductors. In Bi-2212, Bi-2223. Tl-2223 and Tl-2212 compounds, the M₂Ba₂O₅₋₈ (M=Bi, Tl) charge reservoir layers contain trivalent Bi³⁺ and Tl³⁺ ions [7.8.30]. These trivalent ions may localize the carriers at their own sites and can suppress carriers flow to the conducting CuO_2 planes. But the $Cu_{0.5}Tl_{0.5}Ba_2O_{4-\delta}$ charge reservoir layer of our Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-ν}Cd_ν)O_{12-δ} compound has partially insulating/conducting character due to the presence of Cu atoms, which most likely enhances their doping efficiency. The change in the dielectric constants of Cu_{0.5}Tl_{0.5}-Ba₂Ca₃(Cu_{4-v}Cd_v)O_{12- δ} samples is due to the variation of O_{δ} in Cu_{0.5}Tl_{0.5}Ba₂O_{4-δ} charge reservoir layer after oxygen postannealing [15].

The ratio of energy dissipated and energy stored in the material determines the dielectric loss factor ($\tan \delta$) and the absolute value of dielectric loss factor $\tan \delta$ l versus $Log\{f(Hz)\}$ of the oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2-Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ samples are shown in Fig. 4(a–d). The

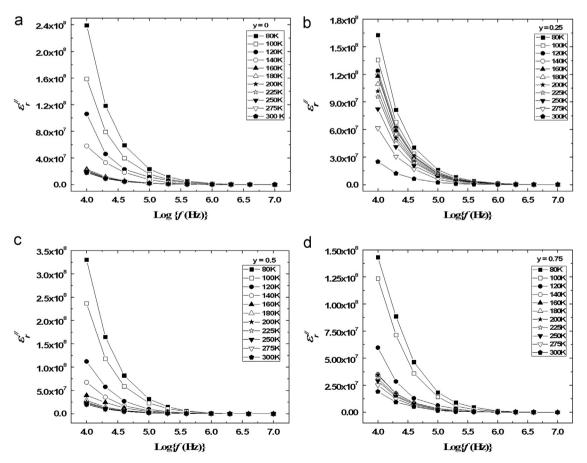


Fig. 3. (a–d). The imaginary part of dielectric (ε''_r) versus log of the frequency $\{\text{Log }f(\text{Hz})\}$ of the external field applied on oxygen post-annealed $\text{Cu}_{0.5}\text{Tl}_{0.5}\text{Ba}_2\text{Ca}_3(\text{Cu}_{4-y}\text{Cd}_y)\text{O}_{12-\delta}$ (a) y=0, (b) y=0.25, (c) y=0.5, and (d) y=0.75, superconductor samples.

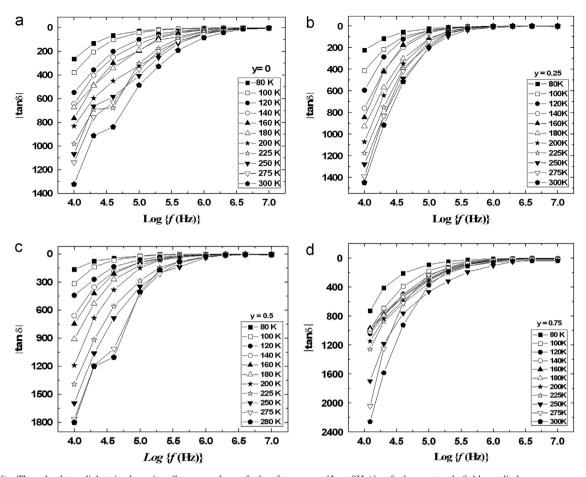


Fig. 4. (a–d). The absolute dielectric loss $\tan \delta l$ versus log of the frequency $\{Log_f(Hz)\}$ of the external field applied on oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ (a) y=0, (b) y=0.25, (c) y=0.5, and (d) y=0.75, superconductor samples.

values of $\tan \delta l$ are observed around 265–1323, 223–1449, 165–1802 and 728–2260 at 10 kHz and temperature range from 80 to 300 K for y=0, 0.25, 0.5, and 0.75, respectively, after oxygen post-annealing $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples. The dielectric loss has increased in all $Cu_{0.5}Tl_{0.5}Ba_2-Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples after oxygen post-annealing. The increase in the value of $l \tan \delta l$ is the conjuncture of oxygen intake after post-annealing in oxygen, which may increase the thickness of the charge reservoir layer and hence the dielectric loss.

The ac-conductivity (σ_{ac}) versus Log{f(Hz)} measurements of oxygen post-annealed Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-y}Cd_y)O_{12- δ} samples have been shown in Fig. 5(a-d). The values of σ_{ac} are observed around 133–10, 60–11, 42–14 and 28–11 for y=0, 0.25, 0.5, and 0.75, respectively, at 10 kHz and in the temperature range of 80–300 K after oxygen post-annealing of Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-y}Cd_y)O_{12- δ} samples. The value of σ_{ac} was decreased significantly in all Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-y}Cd_y) O_{12- δ} samples after oxygen post-annealing. The density of the phonon population is suppressed due the damped harmonic oscillations of heavier Cd atoms in the CuO₂ planes, which reduces the magnitude of superconductivity in the final compound [31]. But after oxygen post-annealing the value of T_c (R=0) is increased, which is most likely due to

optimization of oxygen in $\text{Cu}_{0.5}\text{Tl}_{0.5}\text{Ba}_2\text{O}_{4-\delta}$ charge reservoir layer that optimizes the density of carriers in the conducting planes [15].

These materials with giant dielectric constant are playing an increasingly significant role in the race of microelectronics because of the desire for small and more robust devices such as capacitors and memory devices [32]. However, the main obstacle for potential applications as dielectrics is the high conductivity of cuprates, which results in high dissipation loss. To increase the dielectric constant and to reduce the dissipation loss in superconductors is an open question for all the scientific community.

4. Conclusion

We have investigated the effects of oxygen post-annealing on the superconductivity and dielectric properties of $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ superconductor samples. The negative capacitance (NC) has been observed in all oxygen post-annealed $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-y}Cd_y)O_{12-\delta}$ samples. The change in the dielectric parameters after oxygen post-annealing is the evidence of the variation in O_{δ} oxygen of $Cu_{0.5}Tl_{0.5}Ba_2O_{4-\delta}$ charge reservoir layer that controls the flow of carriers toward the conducting CuO_2 planes. It is observed

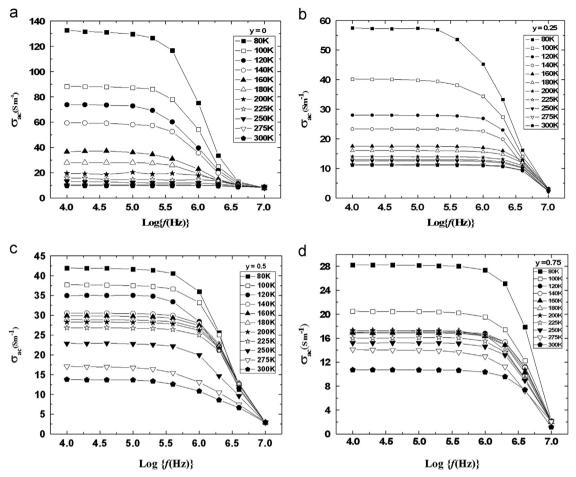


Fig. 5. (a–d). The ac-conductivity versus log of the frequency $\{\text{Log } f(\text{Hz})\}$ of the external field applied on oxygen post-annealed $\text{Cu}_{0.5}\text{Tl}_{0.5}\text{Ba}_2\text{Ca}_3(\text{Cu}_{4-y}\text{Cd}_y)\text{O}_{12-\delta}$ (a) y=0, (b) y=0.25, (c) y=0.5, and (d) y=0.75, superconductor samples.

that the dielectric properties $(\varepsilon'_r, \varepsilon''_r, \text{ and } \sigma_{ac})$ are decreased, while $|\tan \delta|$ is increased in $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ samples after oxygen post-annealing. The variation in O_{δ} of Cu_{0.5}Tl_{0.5}Ba₂O_{4-\delta} charge reservoir layer after oxygen postannealing changes the dielectric properties of Cu_{0.5}Tl_{0.5}Ba₂- $Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ samples. The large intake of oxygen is reflected from appreciable change in the dielectric properties of $Cu_{0.5}Tl_{0.5}Ba_2Ca_3(Cu_{4-\nu}Cd_{\nu})O_{12-\delta}$ samples after oxygen postannealing. At normal state temperature (300 K), the thermal agitation is higher and polarizability is lower resulting in the decrease of dielectric constant and increase in dielectric loss factor. A strong dielectric dispersion is observed at low frequencies and temperatures and the reverse effects are observed on dielectric properties at higher frequencies and temperatures. The observed variations in the dielectric parameters of Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4-ν}Cd_ν)O_{12-δ} samples strongly depend on the working (or operating) temperature and frequency of the external applied ac-field. The dielectric behavior of the materials can be regulated by the oxygen post-annealing as a result of which the variation of oxygen content in the structure of the materials controls the carrier's density. Therefore, the oxygen post-annealing can be used as a tuning technique of these frequency and temperature dependent dielectric properties of Cu_{0.5}Tl_{0.5}Ba₂Ca₃(Cu_{4z-ν}Cd_ν)O_{12-δ} superconducting samples. On the other hand, this oxygen postannealing may affect the inter-grain connectivity of the material.

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