

Eu³⁺ luminescence in novel garnet-type Li₅La₃Nb₂O₁₂ ceramics

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Abstract

A novel Li₅La₃Nb₂O₁₂:Eu³⁺ phosphor was synthesized by the solid-state reaction at high temperature. The phase purity of products was checked by XRD. The emission spectrum is independent of the excitation wavelength and the ⁵D₀–⁷F₂ transition is the dominant one. In the excitation spectrum, there are some sharp peaks originating from the f–f transitions of Eu³⁺ and a broad band peaking at about 280 nm. The concentration dependence of emission intensity and decay curves of ⁵D₀–⁷F₂ transition were investigated. The temperature dependence of emission intensity and decay curves were also investigated in detail.

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1. Introduction

Rare earth ions activated luminescence materials have grown considerably in recent years due to their excellent optical properties and their wide applications in the fields of lighting, display, medical diagnosis and so on [1–7]. It is well established that Eu³⁺ activated phosphors emit in the 580–625 nm spectral range depending on the local environment of Eu³⁺. Recently, Eu³⁺ activated luminescent materials have been paid more attention as interesting candidates for potential red-emitting phosphors in white light-emitting diodes (white LEDs) [8]. White LEDs using blue chip combined with yellow phosphor (Y₃Al₅O₁₂:Ce³⁺), although this phosphor exhibits high luminescence efficiency and chemical stability, show a lower color rendering index (CRI) due to the lack of red color contribution. In order to overcome this shortcoming, red phosphor is added for white LEDs having high CRI [9–11]. So it is important for developing new red phosphors to promote the development of high quality white LEDs.

Garnet-type compounds are a large family and their luminescence properties have been investigated widely [12–15]. Among them, YAG:Ce is one of the most famous phosphors [16]. Li₅La₃Nb₂O₁₂ is also a compound having

garnet structure [17]. Many authors have reported the mechanism of Li ion transport in this crystal because of its potential application in Li-ion batteries. To the best of our knowledge, there are no reports concerning the Eu³⁺ luminescence in Li₅La₃Nb₂O₁₂ host. For the synthesis of rare earth ions activated luminescence materials, substitution occurs readily when two or more elements, with similar ionic radii and the same charge can take each other place within a crystal structure without substantially altering the crystal structure. In this case, La³⁺ is the component element in this host. So the La³⁺ ions could be very easy to be substituted by other rare earth ions due to similar ionic radii and the same charge. Hence we think it is worthwhile to investigate the luminescence properties of Eu³⁺ in Li₅La₃Nb₂O₁₂ host. In this paper, we synthesized a series of Li₅La₃Nb₂O₁₂:Eu³⁺ phosphors with different Eu³⁺ concentrations. Their luminescence properties were characterized by emission spectrum (room and high temperatures), excitation spectrum and luminescence decay behavior in detail.

2. Experimental

Li₅(La_{3–x}Eu_x)Nb₂O₁₂ phosphors were prepared by conventional solid-state reaction at high temperature. Raw materials Li₂CO₃ (excess 20 mol% to compensate the

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evaporation), La_2O_3 , Eu_2O_3 , and Nb_2O_5 (all materials are bought from Aldrich, 99.99%) were thoroughly mixed, and sintered in air at 900 °C for 24 h. X-ray powder diffraction (Beijing Puxi, XD-2, Cu $K\alpha$, $\lambda=1.5405 \text{ \AA}$) was used to identify the structure of the final products. Excitation and emission spectra at room temperature were carried out by a fluorescence meter (Shimadzu, RF 5301PC) with a 150 W Xe lamp as an excitation source. The emission spectra and decay curves at the temperature range between 297 and 503 K were recorded by a 500 MHz digital storage oscilloscope (LeCroy 9350 A) in which the signal was fed from PMT. The samples were excited by 266 nm UV laser. The setup was equipped with a home-made heating cell connected to the temperature controller.

3. Results and discussion

As one of the garnet-type compounds, $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ has a cubic unit cell with space group $I2_13$ (no. 199); the lattice parameters are $a=12.797 \text{ \AA}$ and $V=2095.68 \text{ \AA}^3$ [17]. In order to check the phase purity of the samples, XRD measurement was carried out for all samples. The selected XRD patterns of $\text{Li}_5(\text{La}_{3-x}\text{Eu}_x)\text{Nb}_2\text{O}_{12}$ samples ($x=0.02, 0.15, 0.30, 0.50$ and 1.0) are shown in Fig. 1. It is found that all observed diffraction peaks can be indexed to the standard date of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$. No other impurity phases can be detected when Eu^{3+} ions are doped into the host lattice, indicating that the samples are of single phase. However, the XRD patterns are not consistent with the standard diffraction peaks of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ when the value of x is greater than 1.0. The heavy dopant products should have different crystal structures, which could be attributed to the radius difference of Eu^{3+} and La^{3+} ions. The crystal structure of these samples is under investigation.

The Eu^{3+} activated $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ samples emit intense red luminescence under UV excitation (254 nm). It is found that the emission spectrum is independent of the excitation wavelength except for the emission intensity.

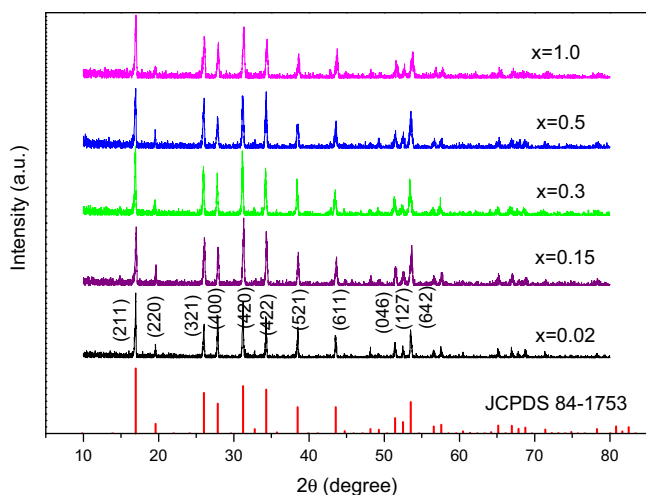


Fig. 1. XRD patterns of $\text{Li}_5(\text{La}_{3-x}\text{Eu}_x)\text{Nb}_2\text{O}_{12}$ samples. Standard data for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ (JCPDS card no. 84–1753) is shown as a reference.

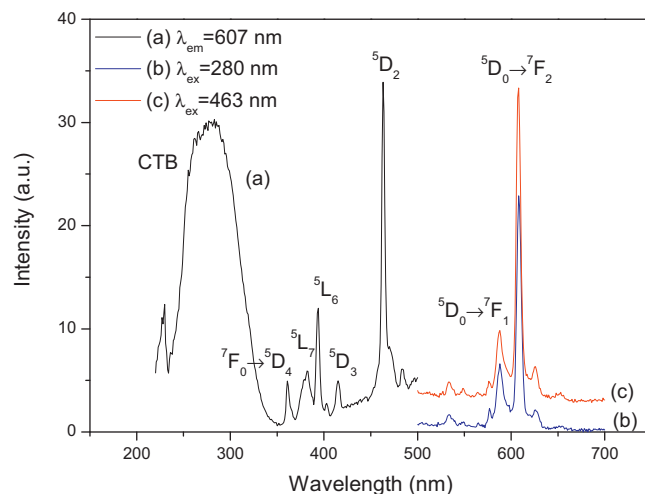


Fig. 2. Emission ($\lambda_{\text{ex}}=280, 463 \text{ nm}$) and excitation ($\lambda_{\text{em}}=607 \text{ nm}$) spectra of $\text{Li}_5(\text{La}_{2.70}\text{Eu}_{0.30})\text{Nb}_2\text{O}_{12}$ phosphor.

Fig. 2 shows representative emission spectra of $\text{Li}_5(\text{La}_{2.70}\text{Eu}_{0.30})\text{Nb}_2\text{O}_{12}$. The research of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ crystal structure indicated that the lanthanum atom is surrounded by four O atoms and forms a contracted tetrahedron [17]. It can be seen from Fig. 2 that the $^5\text{D}_0\text{--}^7\text{F}_2$ transition (electric dipole transition) is the dominant one, which reveals that Eu^{3+} ions are located in the asymmetry sites [18]. So we think that the doped Eu^{3+} ions should substitute the La^{3+} sites. The CIE (Commission Internationale de l'Eclairage) chromaticity coordinates are calculated to be $x=0.53$ and $y=0.45$. The excitation spectrum of the Eu^{3+} emission (607 nm) consists of well-known Eu^{3+} f–f transition lines and a broad band peak at about 280 nm. Usually, for the Eu^{3+} -doped samples, the position of charge transfer band (CTB) can be calculated with the help of an empirical formula given by Reisfeld and Jørgensen [19]

$$E_{\text{ct}}(\text{cm}^{-1}) = [\chi_{\text{opt}}(\text{X}) - \chi_{\text{uncorr}}(\text{M})] \times 30,000 \text{ cm}^{-1} \quad (1)$$

where E_{ct} is the position of the CTB in cm^{-1} , $\chi_{\text{opt}}(\text{X})$ is the optical electronegativity of the anion, and $\chi_{\text{uncorr}}(\text{M})$ is that of the central metal ion. Using the Pauling scale for electronegativity, namely $\chi_{\text{opt}}(\text{O})=3.44$, and $\chi_{\text{uncorr}}(\text{Eu}^{3+})=1.74$ [20], the CTB of Eu–O can be estimated to be $51,000 \text{ cm}^{-1}$ (196 nm). However, this calculated value shows large difference with the experimental data (280 nm). There are two possible reasons which can be applied to explain this red shift [21]. One is that in the Sc–Lu–Y–Gd–La series, the position of Eu–O CTB in La^{3+} host has minimum E_{ct} . Another is that in Eu^{3+} doped composite oxides (Eu–O–M, M is metal element), the energy of CTB depends not only on the coordination atom O, but also on the metal atom M. $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ is a composite oxide. In this case, the electronegativity of Nb is small and this weakens the interaction between M and O, which is beneficial for transferring the electron from O to Eu and decreases the E_{ct} . In addition, the charge transfer transition of host, namely, an electron from an oxygen ligand to the empty 4d orbital of the central

niobium atoms in NbO groups (Nb→O) could also happen in this region. In the crystal structure of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$, Nb atom is surrounded by six oxygen atoms and forms NbO_6 octahedron. Many authors report that NbO_6 group has absorption band at about 300 nm [22–24]. So we regard that the 280 nm absorption band in the excitation spectrum could be attributed to the overlap of CTB of Eu–O and NbO_6 group absorption. The 463 nm excitation peak is strong and matches the emission of blue light emitting diode. So $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}:\text{Eu}^{3+}$ could be a candidate phosphor for phosphor-converted white LEDs.

The concentration dependence of emission intensity of $^5\text{D}_0\text{--}^7\text{F}_2$ transition is shown in Fig. 3. It can be seen that the emission intensity increases with increasing Eu^{3+} concentration, and the optimum value of x in $\text{Li}_5(\text{La}_{3-x}\text{Eu}_x)\text{Nb}_2\text{O}_{12}$ is 0.30. The emission intensity shows similar trend when three different frequencies of excitation light are used (280, 393 and 463 nm). The high doping concentration of activated ions in a phosphor can suffer high power excitation in the application of solid state lighting, which is very useful for devices. Concentration quenching is due to migration of the excitation energy to trap sites [25]. Therefore, concentration quenching often takes place when the average distance between the Eu^{3+} ions is short enough to transfer their energy. So it is possible to estimate the critical distance (R_0 , the distance at which the probability of transfer is equal to the probability of radiative emission) from the concentration quenching data. The critical transfer distance (R_0) can be derived roughly from the following formula proposed by Blasse [26]:

$$R_0 \approx 2 \times \left(\frac{3V}{4\pi x_c N} \right)^{1/3} \quad (2)$$

where x_c is the critical concentration at which the quenching occurs, N is the number of La^{3+} ions in the $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ unit cell and V is the volume of the unit cell. By taking the experimental and analytical values of V ,

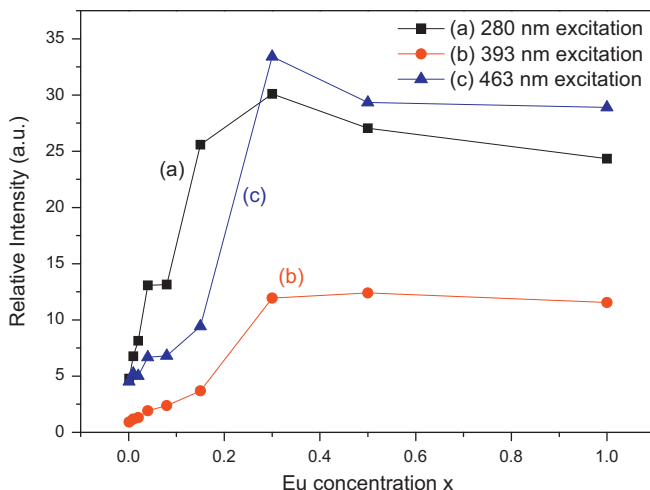


Fig. 3. Emission intensity of $^5\text{D}_0\text{--}^7\text{F}_2$ transition as a function of Eu^{3+} concentration for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}:\text{Eu}^{3+}$ under different excitation lights.

N and x_c (2095.68 \AA^3 , 8, 0.30, respectively), the critical transfer distance of Eu^{3+} in $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}:\text{Eu}^{3+}$ phosphor is calculated to be 12 Å, which indicates that the mechanism of exchange interaction plays no role in energy transfer between Eu^{3+} ions in the $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}:\text{Eu}^{3+}$ phosphors. In general, the exchange interaction is dominant only for short distance (typical critical distances of 5 Å). So the energy transfer in $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}:\text{Eu}^{3+}$ will take place only through electric multipolar interaction.

The decay curves of the 610 nm luminescence of Eu^{3+} in $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}:\text{Eu}^{3+}$ phosphors with different doped concentrations are shown in Fig. 4. It can be seen that all of the curves show exponential decay, and the heavier the doped concentration, the shorter the lifetime. The decay curves can be fitted by single exponential equation: $I(t) = I_0 + A \exp(-t/\tau)$, where I and I_0 is the luminescence intensity, A is a constant, t is the time, τ is the decay time for the exponential components. The τ extracted from the fitted curve is 1.14 ms for $\text{Li}_5(\text{La}_{2.999}\text{Eu}_{0.001})\text{Nb}_2\text{O}_{12}$ and 0.37 ms for $\text{Li}_5(\text{La}_2\text{Eu})\text{Nb}_2\text{O}_{12}$. When the doped concentration increases, the decrease in lifetime can be attributed to the non-radiative relaxation between Eu^{3+} ions. When the doped concentration increases, the sample has more and more quenching centers and the non-radiative rate increases, which causes the lifetime to be shortened [27].

The thermal quenching property of phosphor is one of the important issues to be considered in white LEDs especially in high-power devices. The temperature dependences of PL emission spectra of $\text{Li}_5(\text{La}_{2.70}\text{Eu}_{0.30})\text{Nb}_2\text{O}_{12}$ excited by 266 nm light are shown in Fig. 5. With increase in temperature, the thermal quenching takes places and the emission intensity decreases obviously. The thermal quenching is due to the thermal relaxation through the crossing point between the excited state and the ground state in the configurational coordinate diagram [18]. A non-radiative transition occurs because of the interaction between the electron and thermally active phonon [28].

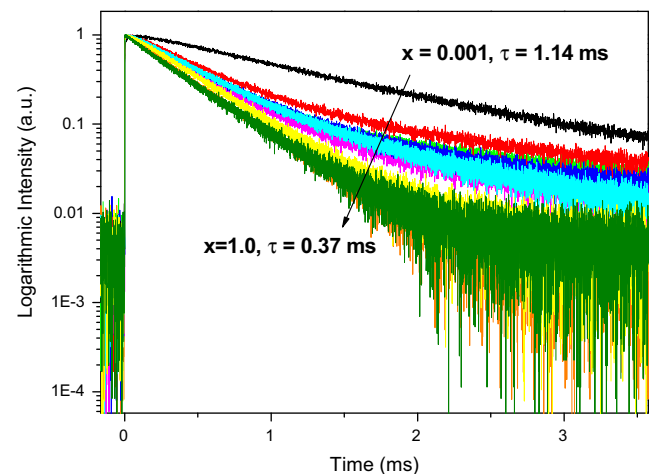


Fig. 4. Decay curves of the 610 nm luminescence of Eu^{3+} as a function of Eu^{3+} concentration in $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}:\text{Eu}^{3+}$ phosphors excited at 266 nm.

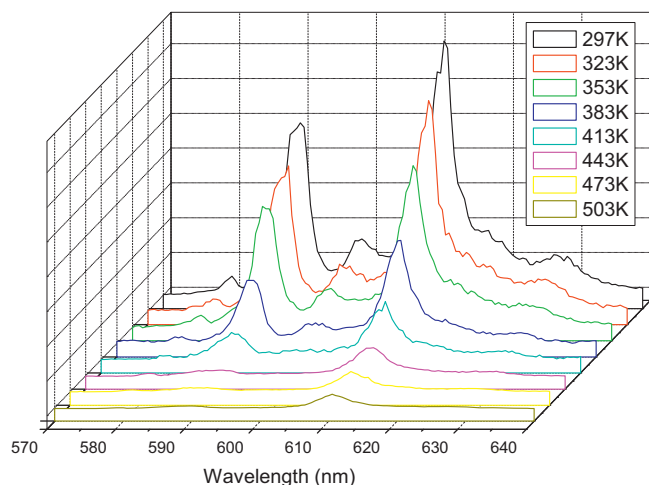


Fig. 5. Emission spectra of the Eu^{3+} as a function of temperature dependence for $\text{Li}_5(\text{La}_{2.70}\text{Eu}_{0.30})\text{Nb}_2\text{O}_{12}$.

The thermal activation energy for the thermal quenching of the Eu^{3+} emission was determined by measuring the temperature dependence of the Eu^{3+} emission intensity. The temperature dependence of the luminescence intensity is described by a modified Arrhenius equation [29]:

$$I(T) = \frac{I_0}{1 + c \exp(-E/kT)} \quad (3)$$

where I_0 is the initial intensity, $I(T)$ is the intensity at a given temperature T , c is a constant, E is the activation energy for thermal quenching and k is the Boltzmann constant. Fig. 6 plots $\ln[(I_0/I) - 1]$ vs. $1/T$. When performing linear fitting, the obtained value of slope is -4121 . According to Eq. (2), the calculated activation energy is about 0.36 eV.

The decay curves of the Eu^{3+} emission (610 nm) as a function of temperature for $\text{Li}_5(\text{La}_{2.70}\text{Eu}_{0.30})\text{Nb}_2\text{O}_{12}$ phosphor excited at 266 nm are shown in Fig. 7. It is found that the lifetimes decrease from 0.54 ms at room temperature to 0.39 ms at 503 K. The reason is that the lifetime of excited states decreases with increase in temperature, and then resulting in a substantial increase in non-radiative relaxation rate (the radiative relaxation rate usually also increases slightly) [30].

4. Conclusions

$\text{Li}_5(\text{La}_{3-x}\text{Eu}_x)\text{Nb}_2\text{O}_{12}$ phosphors were prepared by solid state reaction. XRD patterns reveal that the samples are of single phase when $x \leq 1.0$. All the samples show red luminescence. The $^5\text{D}_0$ – $^7\text{F}_2$ transition is the dominant one indicating that the Eu^{3+} ion occupies a site with a deviation from inversion symmetry. The CIE chromaticity coordinates are $x=0.53$ and $y=0.45$. The optimum concentration of Eu^{3+} ions is 0.30. The critical transfer distance of Eu^{3+} in $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}:\text{Eu}^{3+}$ phosphor is about 12 Å. The excitation spectrum consists of a strong Eu^{3+} f–f excitation peak at 463 nm and a broad band peak at about 280 nm. The decay curves show single exponential decay and that the lifetimes

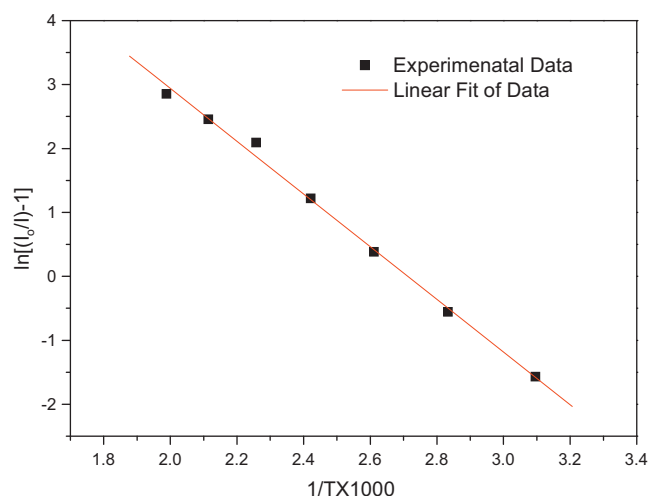


Fig. 6. Activation energy for thermal quenching of $\text{Li}_5(\text{La}_{2.70}\text{Eu}_{0.30})\text{Nb}_2\text{O}_{12}$.

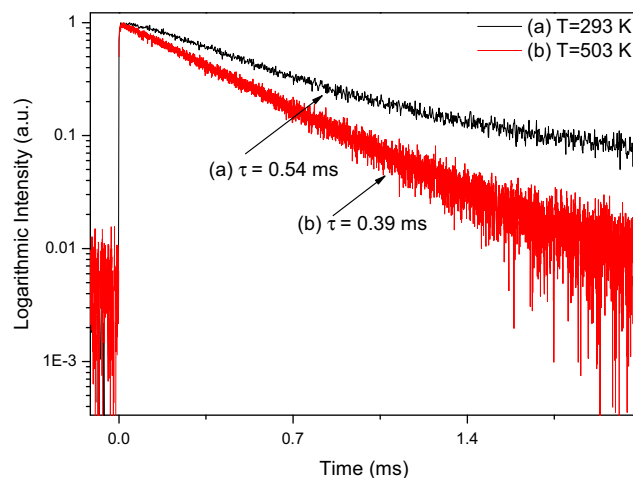


Fig. 7. Decay curves of the Eu^{3+} emission (610 nm) as a function of temperature for $\text{Li}_5(\text{La}_{2.70}\text{Eu}_{0.30})\text{Nb}_2\text{O}_{12}$ phosphor excited at 266 nm.

become shorter with increasing Eu^{3+} concentration. The decrease in lifetime can be attributed to the non-radiative relaxation between Eu^{3+} ions. At high temperature, the thermal quenching takes places and the emission intensity decreases obviously. The thermal quenching can be explained with configurational coordinate diagram. The lifetimes also decrease at high temperature because of a substantial increase in non-radiative relaxation rate.

Acknowledgments

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